

ESTRUCTURA DE LA MATERIA 4

CURSO DE VERANO 2021

CLASE 17

RODOLFO SASSOT

CLASE 17: Neutrinos, CP

Temas: masas de neutrinos, el problema de los neutrinos solares, oscilaciones, CP

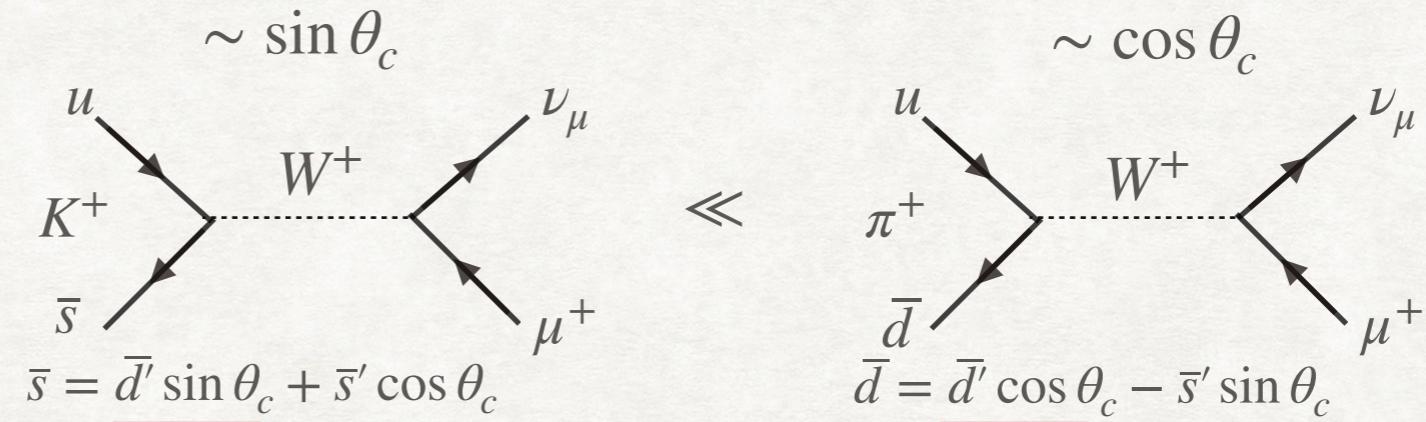
mecanismo de ruptura espontánea genera masas para quarks físicos

(y no para los autoestados débiles d', s', b')

de la clase 14

$$\begin{pmatrix} \nu_e \\ e^- \end{pmatrix} \begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix} \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix} \begin{pmatrix} u \\ d \end{pmatrix} \begin{pmatrix} c \\ s \end{pmatrix} \begin{pmatrix} t \\ b \end{pmatrix}$$

Cabibbo Kobayashi Maskawa



$$\begin{pmatrix} u \\ d' \end{pmatrix} \begin{pmatrix} c \\ s' \end{pmatrix} \begin{pmatrix} t \\ b' \end{pmatrix} \quad \begin{cases} d' \equiv d \cos \theta_c + s \sin \theta_c \\ s' \equiv -d \sin \theta_c + s \cos \theta_c \end{cases} \quad \begin{cases} d \equiv d' \cos \theta_c - s' \sin \theta_c \\ s \equiv d' \sin \theta_c + s' \cos \theta_c \end{cases} \quad \begin{aligned} \theta_c &\simeq 13^\circ \\ \sin^2 \theta_c &\simeq 0.05 \end{aligned}$$

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} U_{dd} & U_{ds} & U_{db} \\ U_{sd} & U_{ss} & U_{sb} \\ U_{bd} & U_{bs} & U_{bb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$\begin{aligned} m_d &= 4.7 {}^{+0.5}_{-0.3} \text{ MeV} \\ m_s &= 95 {}^{+9}_{-6} \text{ MeV} \\ m_b &= 4180 {}^{+40}_{-30} \text{ MeV} \end{aligned}$$

$(\nu_e, \nu_\mu, \nu_\tau)_L$ son autoestados débiles, qué pasa si se generaran masas a (ν_1, ν_2, ν_3) ?

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$(\nu_e, \nu_\mu, \nu_\tau)_L$ son autoestados débiles, qué pasa si se generaran masas a (ν_1, ν_2, ν_3) ?

⇒ consecuencias medibles:

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + h(x) \end{pmatrix} \quad \text{interacción con } h$$

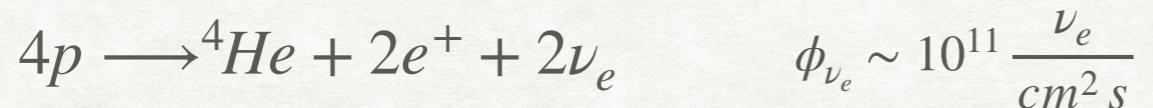
oscilaciones cuánticas $|\Psi(t)\rangle = e^{-i/\hbar E_1 t} |\phi_1\rangle + e^{-i/\hbar E_2 t} |\phi_2\rangle + \dots$

(R. Davies M. Koshiba Nobel 2002, A. McDonald T. Kajita 2015)



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el problema de los neutrinos solares (~1968)

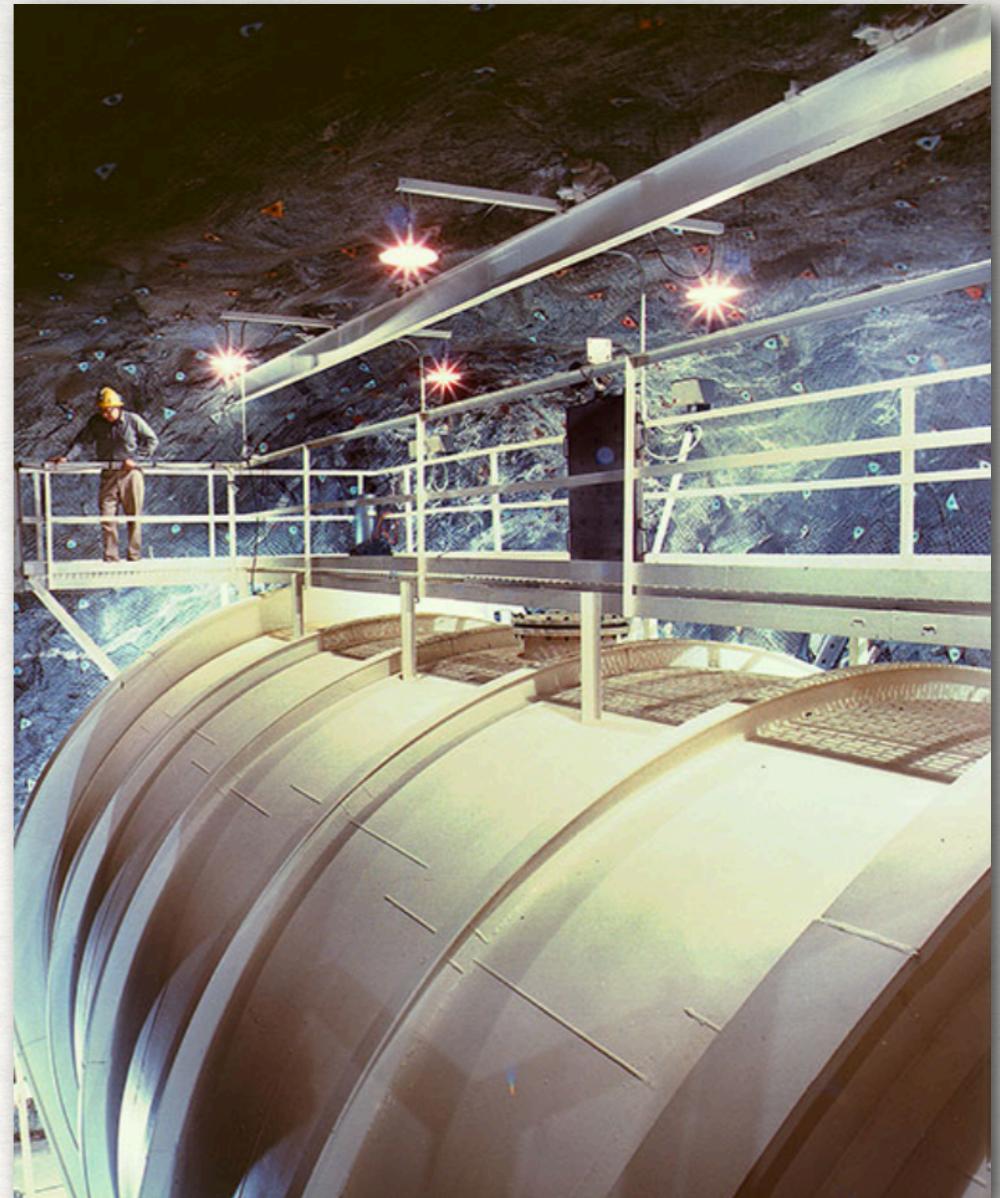


1964 R. Davies, J. Bahcall

40.000 l ${}^{37}\text{Cl}$

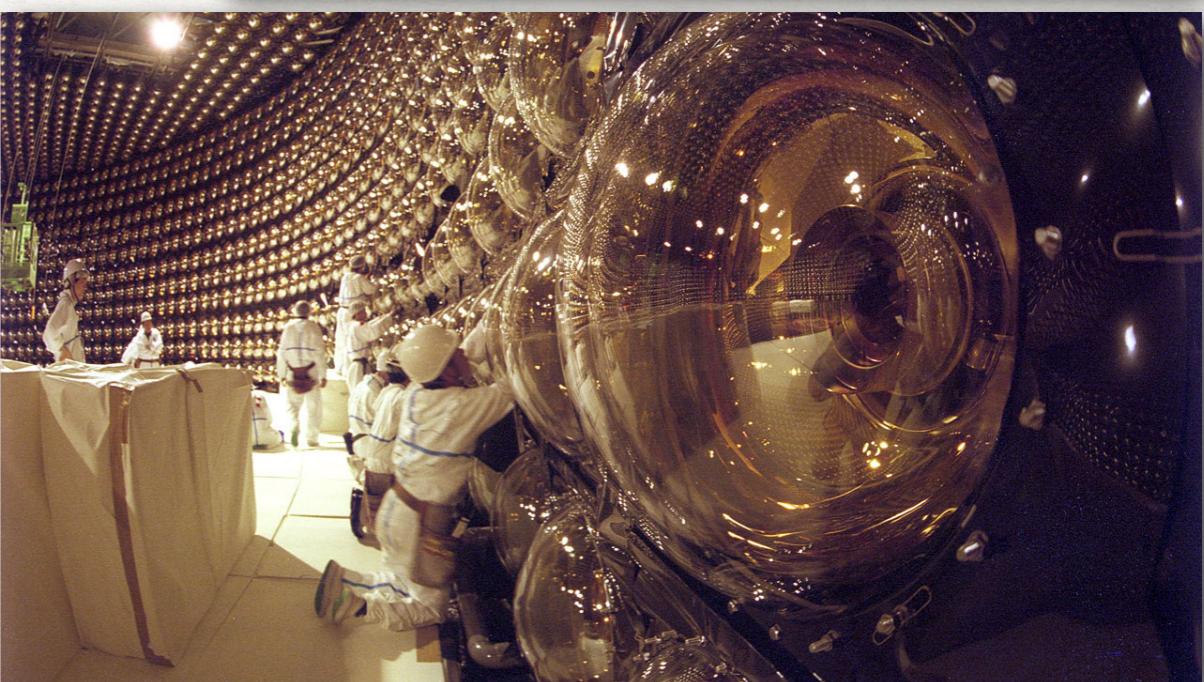


1968 $\sim 1/3$ 'deficit'



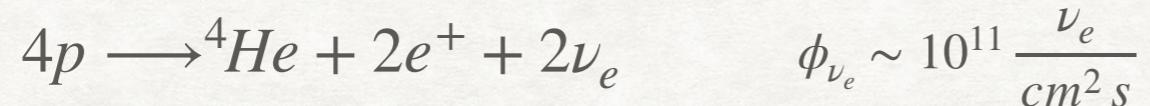
1989 Kamiokande

3.000.000 l H_2O



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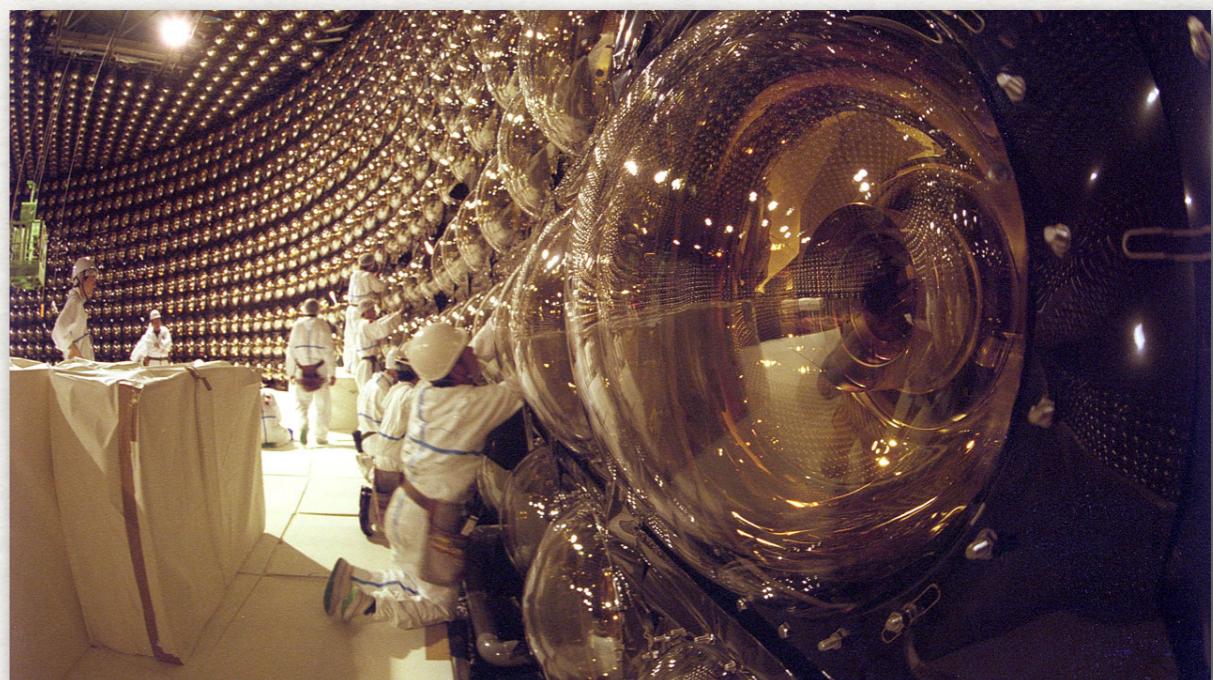
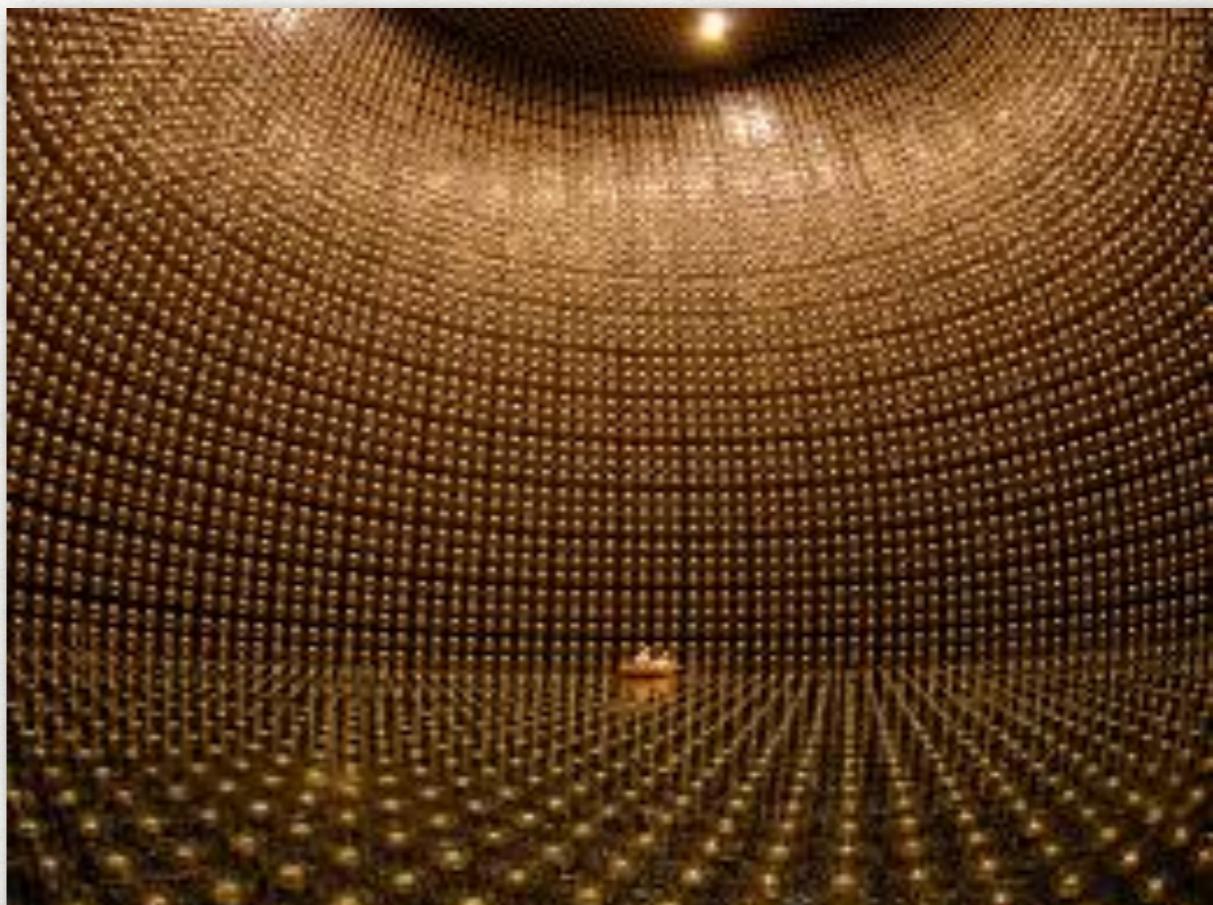
3.000.000 l H_2O

1996 Super-Kamiokande

50.000.000 l H_2O

2001 Sudbury

1.000.000 l D_2O 'esquizofrenia'



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oscilaciones de neutrinos:

$$\begin{pmatrix} \nu_\mu \\ \nu_e \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$



B. Pontecorvo (1957)

$$\begin{cases} \nu_\mu = \nu_1 \cos \theta + \nu_2 \sin \theta \\ \nu_e = -\nu_1 \sin \theta + \nu_2 \cos \theta \end{cases}$$

$$\begin{cases} \nu_1 = \nu_\mu \cos \theta - \nu_e \sin \theta \\ \nu_2 = \nu_\mu \sin \theta + \nu_e \cos \theta \end{cases}$$

$$\frac{\nu_1(t) = \nu_1(0) e^{-iE_1 t}}{\nu_2(t) = \nu_2(0) e^{-iE_2 t}}$$

si se generan con p fijo

$$E_i \simeq p + \frac{m_i^2}{2p} \quad (E_i \gg m_i)$$

si a $t = 0$ $\nu_\mu(0) = 1$

$$\underline{\nu_1(0) = \nu_\mu(0)\cos \theta} \quad \nu_2(0) = \nu_\mu(0)\sin \theta$$

$$\nu_\mu(t) = \underline{\nu_1(t)} \cos \theta + \nu_2(t) \sin \theta = \nu_\mu(0) \cos^2 \theta e^{-iE_1 t} + \nu_\mu(0) \sin^2 \theta e^{-iE_2 t}$$

$$A_\mu = \frac{\nu_\mu(t)}{\nu_\mu(0)} = \cos^2 \theta e^{-iE_1 t} + \sin^2 \theta e^{-iE_2 t}$$

$$A_\mu A_\mu^* = 1 - \sin^2(2\theta) \sin^2 \left(\frac{E_2 - E_1}{2} t \right)$$

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oscilaciones de neutrinos:

$$A_\mu A_\mu^* = 1 - \sin^2(2\theta) \sin^2 \left(\frac{E_2 - E_1}{2} t \right)$$

$$\begin{aligned} \Delta m^2 &\equiv m_2^2 - m_1^2 \simeq 2p(E_2 - p) - 2p(E_1 - p) \\ &\simeq 2p(E_2 - E_1) \end{aligned}$$

$$E_i \simeq p + \frac{m_i^2}{2p}$$

$$P(\nu_\mu \rightarrow \nu_\mu) = A_\mu A_\mu^* = 1 - \sin^2(2\theta) \sin^2 \left(\frac{\Delta m^2}{4p} t \right) = 1 - \sin^2(2\theta) \sin^2 \left(\frac{\Delta m^2 c^4 L}{4E\hbar c} \right)$$

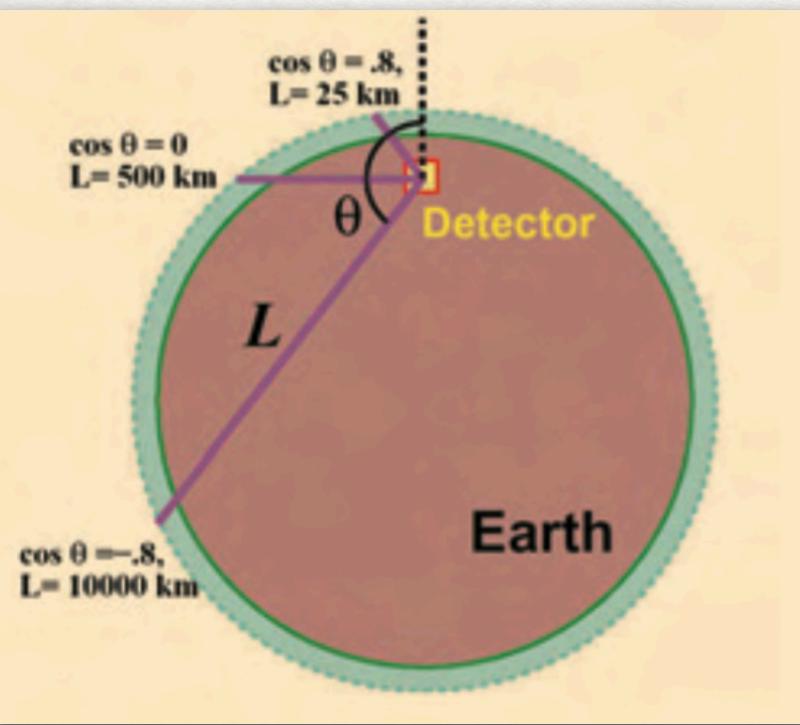
$$E_i \gg m_i \quad v \simeq c \quad E \simeq p \quad t \simeq \frac{L}{c}$$

$$\Delta m_{12}^2 = 7.53 \pm 0.18 \ 10^{-5} \text{ eV}^2$$

$$\Delta m_{23}^2 = 2.546 \pm 0.034 \ 10^{-3} \text{ eV}^2$$

$$\theta_{12} \simeq 17.89^\circ$$

$$\theta_{23} \simeq 9.17^\circ$$



B. Pontecorvo (1957)



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oscilaciones de kaones: violación de CP

$$\Gamma(\pi^+ \rightarrow \mu^+ + \nu_L) \neq \Gamma(\pi^+ \rightarrow \mu^+ + \nu_R) \neq \Gamma(\pi^- \rightarrow \mu^- + \bar{\nu}_R)$$

CP

P C



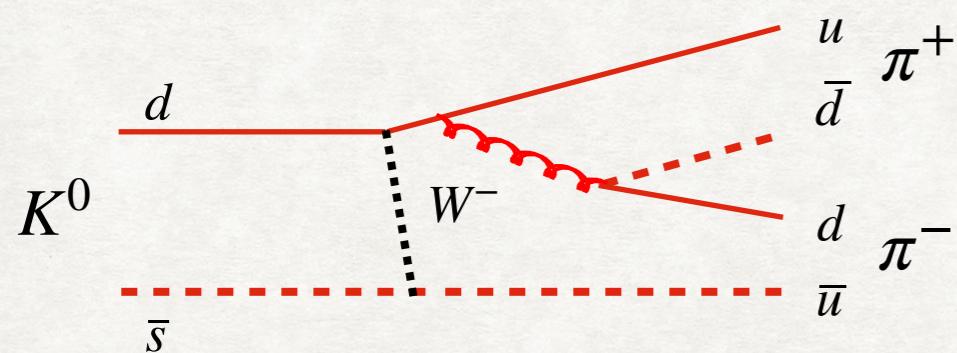
V. Fitch J. Cronin (Nobel 1980)

teorema CPT : invariancia Lorentz + causalidad + QFT implican que CPT es buena simetría

1964 Violación de CP oscilaciones de kaones neutros ($K^0 \equiv d\bar{s}$, $\bar{K}^0 \equiv s\bar{d}$)

2001 en mesones B (SLAC) ($B^0 \equiv d\bar{b}$, $\bar{B}^0 \equiv b\bar{d}$)

2019 en mesones D (LHC) ($D^0 \equiv c\bar{u}$, $\bar{D}^0 \equiv u\bar{c}$)



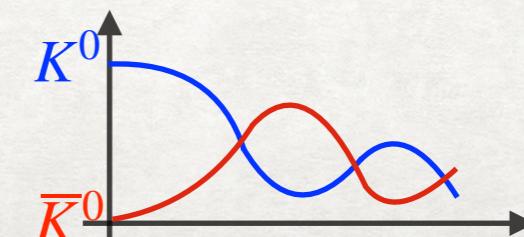
$$CP |K^0\rangle = - |\bar{K}^0\rangle$$

$$|K_L\rangle = \frac{1}{\sqrt{2}} [|K^0\rangle + |\bar{K}^0\rangle] \quad CP |K_L\rangle = - |K_L\rangle$$

$$|K_S\rangle = \frac{1}{\sqrt{2}} [|K^0\rangle - |\bar{K}^0\rangle] \quad CP |K_S\rangle = + |K_S\rangle$$



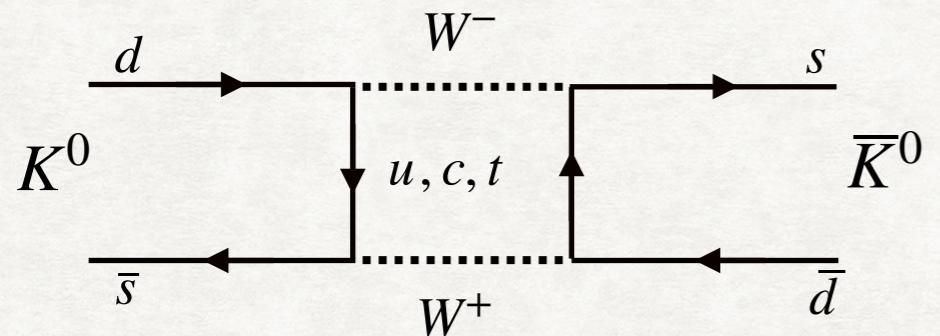
$$|K_0\rangle = \frac{1}{\sqrt{2}} [|K_L\rangle + |K_S\rangle]$$



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oscilaciones de kaones: violación de CP

$$\frac{K_L \rightarrow \pi^0 \pi^0}{K_S \rightarrow \pi^0 \pi^0} = (2.28 \pm 0.02) 10^{-3}$$



$$\Gamma(K^0 \rightarrow \bar{K}^0) \neq \Gamma(\bar{K}^0 \rightarrow K^0)$$

$$\begin{cases} K_L \rightarrow e^+ + \nu_e + \pi^- \\ K_L \rightarrow e^- + \bar{\nu}_e + \pi^+ \end{cases} \quad \text{materia/anti-materia}$$



asimetría bariónica

$$\frac{\Gamma(K_L \rightarrow e^+ + \nu_e + \pi^-) - \Gamma(K_L \rightarrow e^- + \bar{\nu}_e + \pi^+)}{\Gamma(K_L \rightarrow e^+ + \nu_e + \pi^-) + \Gamma(K_L \rightarrow e^- + \bar{\nu}_e + \pi^+)} = (3.32 \pm 0.06) 10^{-3}$$