

ESTRUCTURA DE LA MATERIA 4

CURSO DE VERANO 2021

CLASE 6

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CLASE 6: Fenomenología de Dirac

Temas: Límite no relativista, Fenomenología.

recapitulación:

$$\begin{cases} E \rightarrow i\hbar \frac{\partial}{\partial t} & H = c \vec{\alpha} \cdot \vec{p} + \beta mc^2 & E^2 = c^2 p^2 + m^2 c^4 \\ \vec{p} \rightarrow -i\hbar \vec{\nabla} & \alpha_i^2 = 1 \quad \beta^2 = 1 \quad \{\alpha_i, \alpha_j\} = 2\delta_{ij} \quad \{\alpha_i, \beta\} = 0 \end{cases}$$

$$H\psi = E\psi \quad \psi = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix} \quad \begin{array}{l} 2 \text{ sol. l.i. } E > 0 \\ 2 \text{ sol. l.i. } E < 0 \end{array}$$

$$\rho \equiv \psi^\dagger \psi \geq 0$$

$$\psi^{(1)}(\vec{x}, t) = N \begin{pmatrix} 1 \\ 0 \\ \left(\frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2}\right) 1 \\ 0 \end{pmatrix} e^{-\frac{i}{\hbar} p_\mu x^\mu}$$

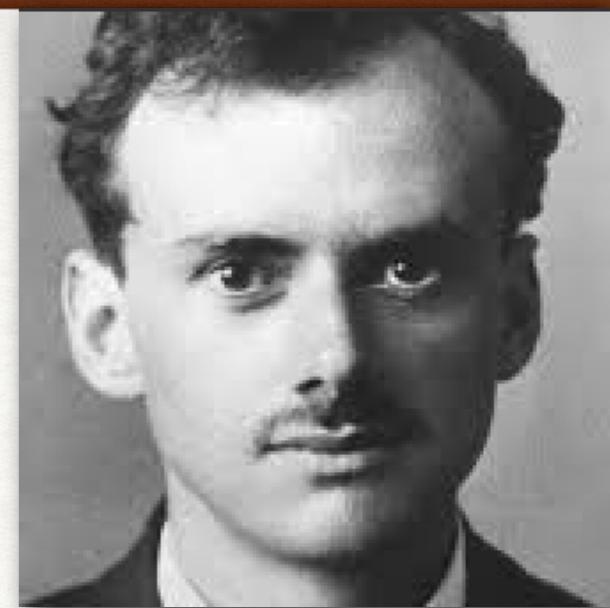
$$\psi^{(2)}(\vec{x}, t) = N \begin{pmatrix} 0 \\ 1 \\ \left(\frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2}\right) 0 \\ 1 \end{pmatrix} e^{-\frac{i}{\hbar} p_\mu x^\mu}$$

$$\Sigma_3 \equiv \begin{pmatrix} \sigma_3 & 0 \\ 0 & \sigma_3 \end{pmatrix} \quad [H, \Sigma_i] \neq 0 \quad \vec{J} \equiv \vec{L} + \frac{\hbar}{2} \vec{\Sigma} \quad [H, J_i] = 0$$



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límite no relativista



$$(c \vec{\alpha} \cdot \vec{p} + \beta mc^2) \psi = E \psi$$

$$\vec{p} \rightarrow \vec{p} - \frac{q}{c} \vec{A} \quad E \rightarrow E - q\phi \quad \psi = \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix}$$

$$\rightarrow \left[c \vec{\alpha} \cdot \left(\vec{p} - \frac{q}{c} \vec{A} \right) + \beta mc^2 \right] \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix} = (\epsilon_{NR} + mc^2 - q\phi) \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix}$$

$\equiv \vec{\pi}$

$$\left[c \begin{pmatrix} 0 & \vec{\sigma} \cdot \vec{\pi} \\ \vec{\sigma} \cdot \vec{\pi} & 0 \end{pmatrix} + \begin{pmatrix} mc^2 & 0 \\ 0 & -mc^2 \end{pmatrix} + q\phi \right] \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix} = (\epsilon_{NR} + mc^2) \begin{pmatrix} \psi_A \\ \psi_B \end{pmatrix}$$

$$c \vec{\sigma} \cdot \vec{\pi} \psi_B + \cancel{mc^2} \psi_A + q\phi \psi_A = (\epsilon_{NR} + \cancel{mc^2}) \psi_A \quad (\epsilon_{NR} - q\phi) \psi_A = c \vec{\sigma} \cdot \vec{\pi} \psi_B$$

$$c \vec{\sigma} \cdot \vec{\pi} \psi_A - mc^2 \psi_B + q\phi \psi_B = (\epsilon_{NR} + mc^2) \psi_B \quad (\epsilon_{NR} + 2mc^2 - q\phi) \psi_B = c \vec{\sigma} \cdot \vec{\pi} \psi_A$$

$$\psi_A = \frac{c \vec{\sigma} \cdot \vec{\pi}}{(\epsilon_{NR} - q\phi)} \psi_B$$

$$\psi_B = \frac{c \vec{\sigma} \cdot \vec{\pi}}{(\epsilon_{NR} + 2mc^2 - q\phi)} \psi_A$$

$$\psi_A = \frac{c \vec{\sigma} \cdot \vec{\pi}}{(\epsilon_{NR} - q\phi)} \frac{c \vec{\sigma} \cdot \vec{\pi}}{(\epsilon_{NR} + 2mc^2 - q\phi)} \psi_A$$

$$\psi_A = \frac{c^2 (\vec{\pi}^2 + i \vec{\sigma} \cdot (\vec{\pi} \times \vec{\pi}))}{(\epsilon_{NR} - q\phi)(\epsilon_{NR} + 2mc^2 - q\phi)} \psi_A \quad \vec{\pi} \times \vec{\pi} \neq 0$$

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$$\begin{aligned}\vec{\pi} \times \vec{\pi} &= (\vec{p} - \frac{q}{c} \vec{A}) \times (\vec{p} - \frac{q}{c} \vec{A}) \\ &= \cancel{\vec{p} \times \vec{p}} - \frac{q}{c} \vec{p} \times \vec{A} - \frac{q}{c} \vec{A} \times \vec{p} + \frac{q^2}{c^2} \vec{A} \times \cancel{\vec{A}} \\ &= \frac{q}{c} i\hbar (\vec{\nabla} \times \vec{A} + \vec{A} \times \vec{\nabla})\end{aligned}$$

$$\begin{aligned}(\vec{\pi} \times \vec{\pi})\psi &= \frac{q}{c} i\hbar (\vec{\nabla} \times \vec{A} \psi + \vec{A} \times \vec{\nabla} \psi) = \frac{q}{c} i\hbar \left[(\vec{\nabla} \times \vec{A})\psi + \vec{\nabla} \psi \times \cancel{\vec{A}} + \vec{A} \times \cancel{\vec{\nabla} \psi} \right] \\ &= \frac{q}{c} i\hbar (\vec{\nabla} \times \vec{A})\psi\end{aligned}$$

$$\psi_A = \frac{c^2 (\pi^2 + i\vec{\sigma} \cdot (\vec{\pi} \times \vec{\pi}))}{(\epsilon_{NR} - q\phi)(\epsilon_{NR} + 2mc^2 - q\phi)} \psi_A = \frac{c^2 \left(\pi^2 + i\vec{\sigma} \cdot (i\hbar \frac{q}{c} (\vec{\nabla} \times \vec{A})) \right)}{(\epsilon_{NR} - q\phi)(\epsilon_{NR} + 2mc^2 - q\phi)} \psi_A$$

$$\psi_A = \frac{c^2 \left((\vec{p} - \frac{q}{c} \vec{A})^2 - \frac{q}{c} \hbar \vec{\sigma} \cdot \vec{B} \right)}{(\epsilon_{NR} - q\phi)(\epsilon_{NR} + 2mc^2 - q\phi)} \psi_A$$

$$\begin{aligned}\psi_A (\epsilon_{NR} - q\phi) &= \frac{1}{2m} \left((\vec{p} - \frac{q}{c})^2 - \frac{q}{c} \hbar \vec{\sigma} \cdot \vec{B} \right) \psi_A \\ &\quad \left[\frac{1}{2m} \left((\vec{p} - \frac{q}{c} \vec{A})^2 - \frac{q}{c} \hbar \vec{\sigma} \cdot \vec{B} \right) + q\phi \right] \psi_A = \epsilon_{NR} \psi_A\end{aligned}$$

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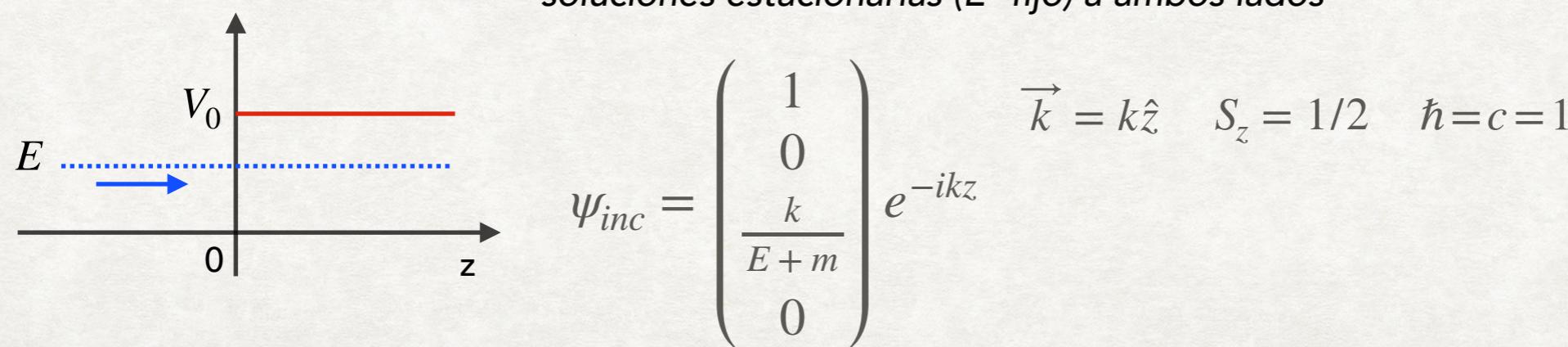
$$\left[\frac{1}{2m} \left((\vec{p} - \frac{\vec{q}}{c})^2 - \frac{q}{c} \hbar \vec{\sigma} \cdot \vec{B} \right) + q\phi \right] \psi_A = \epsilon_{NR} \psi_A \quad \psi_A = \psi_{schroedinger} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\frac{q}{2mc} \hbar \vec{\sigma} \cdot \vec{B} = \frac{gq}{2mc} \vec{S} \cdot \vec{B} \quad g = 2$$

$g = 2.00321930436182(52)$ (NIST 2018)

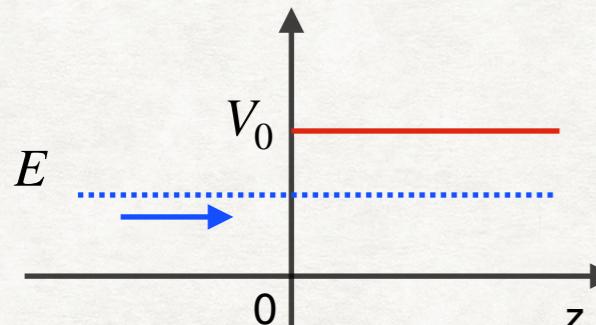
↑
teoría (3 loops)

paradoja de Klein: *cuando las componentes débiles no son tan débiles...*



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paradoja de Klein: cuando las componentes débiles no son tan débiles...



$$\psi_{inc} = \begin{pmatrix} 1 \\ 0 \\ \frac{k}{E+m} \\ 0 \end{pmatrix} e^{-ikz} \quad \vec{k} = k\hat{z} \quad S_z = 1/2 \quad \hbar = c = 1$$

$$\psi_{refl} = a \begin{pmatrix} 1 \\ 0 \\ \frac{-k}{E+m} \\ 0 \end{pmatrix} e^{ikz} + b \begin{pmatrix} 0 \\ 1 \\ 0 \\ \frac{k}{E+m} \end{pmatrix} e^{ikz}$$

$$E = \sqrt{c^2 q^2 + m^2 c^4} + V_0$$

$$E - V_0 = \sqrt{c^2 q^2 + m^2 c^4}$$

$$cq = \sqrt{(E - V_0)^2 - m^2 c^4}$$

$$q = \sqrt{(E - V_0)^2 - m^2}$$

continuidad de ψ en $z=0$ $\psi_{inc}(z=0) + \psi_{refl}(z=0) = \psi_{trans}(z=0)$

$$1 + a = c$$

$$b = d$$

$$k(1 - a)/(E + m) = cq/(E - V_0 + m)$$

$$bk/(E + m) = -dq/(E - V_0 + m)$$

$$\rightarrow (1 - a) = cr$$

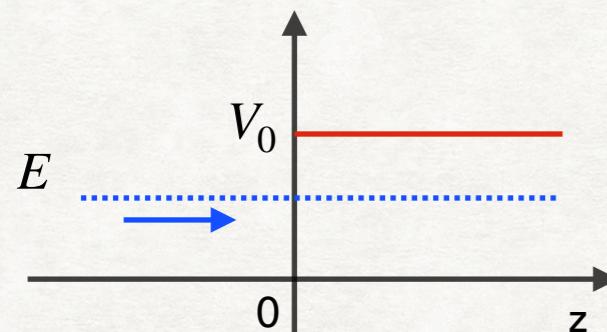
$$\Rightarrow b = d = 0$$

$$r \equiv \frac{q}{k} \frac{E + m}{E - V_0 + m}$$



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paradoja de Klein: cuando las componentes débiles no son tan débiles...



$$1 + a = c$$

$$(1 - a) = cr$$

$$r \equiv \frac{q}{k} \frac{E + m}{E - V_0 + m}$$

$$q = \sqrt{(E - V_0)^2 - m^2}$$



corrientes: $J^\mu \equiv \bar{\psi} \gamma^\mu \psi$ $J^3 \equiv \bar{\psi} \gamma^3 \psi = \psi^\dagger \gamma^0 \gamma^3 \psi = \psi^\dagger \beta \beta \alpha_3 \psi = \psi^\dagger \alpha_3 \psi$

$$J_{inc}^3 = \left(1, 0, \frac{k}{E+m}, 0 \right) \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ \frac{k}{E+m} \\ 0 \end{pmatrix} = \frac{2k}{E+m}$$

$$J_{refl}^3 = \frac{2a^2 k}{E+m}$$

$$J_{trans}^3 = \frac{2c^2 q}{E - V_0 + m}$$

$$\alpha_i = \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix}$$

$$J_{inc}^3 = J_{refl}^3 + J_{trans}^3$$

$$T \equiv \frac{J_{trans}^3}{J_{inc}^3} = \frac{4r}{(1+r)^2} \quad R \equiv \frac{J_{refl}^3}{J_{inc}^3} = \frac{(1-r)^2}{(1+r)^2}$$

a) $V_0 < E + m$ $r > 0$ $T > 0$ *hay transmisión*

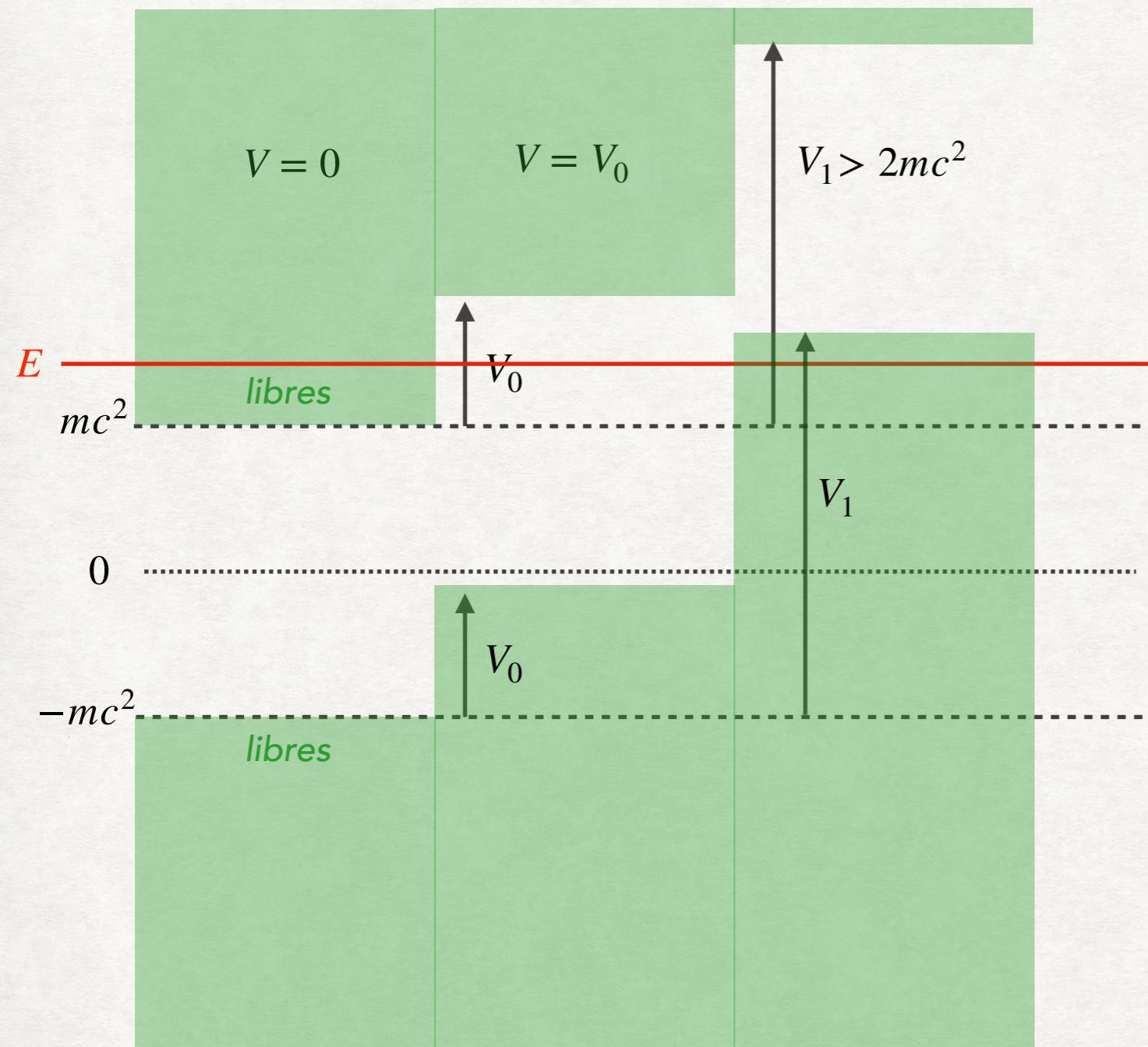
$R < 1$ *se reflejan menos que las incidentes*

b) $V_0 > E + m$ $r < 0$ $T < 0$ *transmitidas en dirección opuesta?*

$R > 1$ *más partículas reflejadas que las incidentes?*

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paradoja de Klein: *cuando las componentes débiles no son tan débiles...*



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trayectoria de una partícula de Dirac: *zitterbewegung*

velocidad de una partícula de Dirac:

$$\dot{x}_k \equiv \frac{i}{\hbar}[H, x_k] = \frac{ic}{\hbar}[\alpha_j p_j, x_k] = c\alpha_k$$

$$H = c \overrightarrow{\alpha} \cdot \overrightarrow{p} + \beta mc^2 \quad [x, p] = i\hbar$$

$$[H, p_k] = 0 \quad [H, \alpha_k] \neq 0$$

aceleración de una partícula de Dirac:

$$\dot{\alpha}_k \equiv \frac{i}{\hbar}[H, \alpha_k] = \frac{i}{\hbar}(-2\alpha_k H + 2cp_k) \neq 0$$

$$\alpha_k(t) = \frac{cp_k}{H} + (\alpha_k(0) - \frac{cp_k}{H})e^{-2iHt/\hbar} \quad (\nu \sim 1.5 \cdot 10^{-21} s^{-1})$$

$$x_k(t) = x_k(0) + \frac{c^2 p_k}{H} t + \frac{ic\hbar}{2H}(\alpha_k(0) - \frac{cp_k}{H})e^{-2iHt/\hbar} \quad (\Delta x \sim 3.9 \cdot 10^{-11} cm)$$



nunca va estar en autotestado de velocidad
(salvo que $m = 0$)

además, $[\alpha_i, \alpha_j] \neq 0 \dots$

para una partícula libre, la velocidad no es constante de movimiento (salvo que $m = 0$)

partícula libre...

